(LATEX: ms cjaa.tex; printed on April 8, 2008; 3:40)

# The debeamed luminosity, sychrotron peak frequency and black hole mass of BL Lac objects

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**Abstract** We estimate the intrinsic luminosities and synchrotron peak frequency using the derived Doppler factor for a sample of 170 BL Lac objects, of which the synchrotron peak frequency are derived by fitting the SED constructed with the collected multiband data from literatures. We find that the debeamed radio and optical core luminosities follow the same correlation found for FR I radio galaxies, which is in support of the unification of the BL Lac objects and the FR I galaxies based on orientation. For the debeamed luminosity at synchrotron peak frequency, we find a significant positive correlation between the luminosity and intrinsic synchrotron peak frequency. This implies that the more powerful sources may have the majority of jet emission at higher frequency. At synchrotron peak frequency, the intrinsic luminosity and black hole mass show strong positive correlation, while mild correlation is found in the case of jet power, indicating that the more powerful sources may have heavier black hole.

**Key words:** black hole physics — BL Lacertae objects: general — galaxies: active — galaxies: jets — galaxies: nuclei

## 1 INTRODUCTION

BL Lac objects are one of the most extreme classes of active galactic nuclei (AGN) with no emission lines or less of them, but they have a strong highly variable and polarized non-thermal continuum emission ranging from radio to γ-ray bands. Traditionally, they have been divided into radio-selected and X-ray-selected BL Lac objects (RBLs and XBLs, respectively) based on the surveys they were primarily discovered. In recent years, the dichotomy has been replaced by a more physically meaningful classification based upon the overall spectral energy distribution (SED) of the object, namely, low frequency peaked BL Lac objects (LBL), intermediate objects (IBL) and high frequency peaked BL Lac objects (HBL) (Padovani & Giommi, 1995). Generally, RBLs tend to be LBLs,

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and XBLs tend to be HBLs (see review by Urry & Padovani (1995)). The synchrotron peak frequency of RBLs is usually in the radio/IR band, while UV/X-ray band for XBLs (Giommi et al., 1995). The recent studies on the SED of large samples of BL Lac objects shown that BL Lac objects most likely form one class with a continuous distribution of synchrotron emission peak energies, while RBLs and XBLs represent the opposite ends of the continuum (Nieppola et al., 2006).

Fossati et al. (1998) and Ghisellini et al. (1998) have proposed the so-called 'blazar sequence' with plots of various powers vs. the synchrotron peak frequency  $v_{\text{peak}}$  for a sample of blazars containing flat-spectrum radio quasars (FSRQ) and BL Lac objects, of which an anti-correlation was apparent, with the most powerful sources having relatively small synchrotron peak frequencies and the least powerful ones having the highest  $\nu_{\text{peak}}$ . The theoretical interpretation to these anti-correlations was proposed by Ghisellini et al. (1998): the more powerful sources suffered a larger probability of losing energy and therefore are subjected to more cooling and therefore translates into a lower value of  $\nu_{\rm peak}$ . However, the recent studies have largely changed this scenario. Recently, Nieppola et al. (2006) proposed that there are anti-correlations between the luminosity at 5 GHz, 37 GHz, and 5500 A and  $v_{\text{peak}}$  for a large sample of BL Lac objects, while based on all the correlations with  $\nu_{\text{peak}}$ , they concluded that the blazar sequence scenario is not valid. Through estimating the Doppler factor for a sample of 170 BL Lac objects selected from Nieppola et al. (2006), Wu et al. (2007) found a significant anti-correlation between the total 408 MHz luminosity, which is supposed to be intrinsic, and the intrinsic  $v_{peak}$ , with a large scatter which however undermines the 'blazar sequence'. Moreover, the evidence of low-power LBLs has been recently discovered (Padonavi et al. 2003; Caccianiga & Marchã 2004; Antón & Browne 2005), and the possible discovery of high luminosity HBLs is also reported in the Sedentary survey (Giommi et al., 2005). In addition, the high-power-high- $\nu_{\text{peak}}$  FSRQs were found in the Deep X-ray Radio Blazar Survey (DXRBS), which is both X-ray and radio-selected, though they do not reach the extreme  $\nu_{peak}$  values of HBLs. From all these discoveries, it seems that the blazar sequence in its simplest form cannot be valid (Padovani 2006). By now, what are the main physical parameters that determines the shift of the synchrotron peak frequency of BL Lac objects seems still unclear.

According to the unified scheme for AGNs, BL Lac objects are thought to be low-luminosity radio galaxies viewed at relatively small angles to the line of sight and the relativistic beaming has an enormous effects on the observed luminosities. Generally, the intrinsic (debeamed) luminosity emitted from jet in BL Lac objects can be estimated using the Doppler factor and assuming the jet geometry. Using the empirical relation of Giovannini et al. (2001), Wu et al. (2007) estimated the Doppler factor for a sample of 170 BL Lac objects with the observed radio core observed either with VLA or MERLIN. In this paper, we will correct the beaming effect in the multi-band luminosity and the synchrotron peak frequency directly using the derived Doppler factor of Wu et al. (2007) for the same sample. Our aim is mainly to investigate the correlation between the luminosity at synchrotron peak frequency and the synchrotron peak frequency after eliminating the beaming effect. The layout of this paper is as follows. The sample is given in Sect. 2. In Sect. 3, the various correlation analysis are presented, which involves the multi-band luminosity, the synchrotron peak frequency, and the black hole mass as well. The discussions are shown in Sect. 4, and the results are summarized in Sect. 5. Throughout the work, we adopt the following values for the cosmological parameters:  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ,  $\Omega_{\text{M}} = 0.3$ , and  $\Omega_{\Lambda} = 0.7$ , except an otherwise stated. The spectral index  $\alpha$  is defined as  $f_{\nu} \propto \nu^{-\alpha}$ .

### 2 THE SAMPLE

Nieppola et al. (2006) presented a large sample of BL Lac objects from the Metsähovi radio observatory BL Lac sample consisting of 381 objects selected from the Veron-Cetty & Veron BL Lac Catalogue (Veron-Cetty & Veron 2000), and 17 objects from the literature, of which many sources are from the well-known BL Lac samples like 1Jy, S4, S5, Einstein Medium Sensitivity Survey (EMSS), Einstein Slew Survey, and DXRBS. The authors argued that this sample is supposed to have no selection criteria (other than declination) in addition to the ones in the original surveys. From this sample, Wu et al. (2007) estimated the Doppler factor for a sample of 170 BL Lac objects using the radio core emission from either VLA or MERLIN observations. The sample investigated in this paper is same as that in Wu et al. (2007), i.e. 170 BL Lac objects with estimated  $\delta$ . The Doppler factor was estimated using the correlation found by Giovannini et al. (2001) for radio galaxies,

$$\log P_{ci5} = 0.62 \log P_t + 8.41 \tag{1}$$

where  $P_{ci5}$  is the intrinsic core 5 GHz radio power derived assuming  $\Gamma=5$  (see Giovannini et al. (2001) for details), and  $P_t$  is the total radio power at 408 MHz. From the measured total 408 MHz radio power, we can obtain the intrinsic core radio power. Thus, the Doppler factor can be got from  $P_{co5}=P_{ci5}\delta^{2+\alpha}$  (corresponding to a continuous jet), assuming  $\alpha=0$ , where  $P_{co5}$  is the observed 5 GHz core luminosity (see Wu et al., 2007, for details).

To construct the SED for all 170 BL Lac objects, we collected the radio, optical and X-ray data from NASA/IPAC Extragalactic Database (NED), the Astrophysical Catalogues Support System (CATs) maintained by the Special Astrophysical Observatory, Russia, and the catalogue of integral blazar working group (IBWG)<sup>1</sup>. The R band core luminosity were mainly obtained from Urry et al. (2000), O'Dowd & Urry (2005), and Nilsson et al. (2003). The X-ray data were collected from Laurent-Muehleisen et al. (1999), Donato et al. (2001), Rector & Stocke (2000), Reich et al. (2000), Turriziani et al. (2007), and RBSC-NVSS (2000). The Galactic extinction in optical band were corrected using the value from NED. Following Nieppola et al. (2006), the SED of each source were constructed in the log  $\nu$  – log  $\nu$ F $_{\nu}$  representation in the observer's frame, based on the multi-frequency data. The synchrotron component of the SED was fitted with a parabolic function  $y = Ax^2 + Bx + C$ , therefore, the synchrotron peak frequency in the rest frame was obtained as  $\nu_{\rm peak} = -B(1+z)/2A$ , in which z is redshift.

In principle, whether or not to include X-ray data in the fit can be decided based on the X-ray spectral index fitted from high quality X-ray spectra. A X-ray spectral index of  $\alpha > 1$  means a synchrotron origin of X-ray emission, while  $\alpha < 1$  indicates an inverse Compton process origin, which implies that the X-ray data should be excluded in the fit. Unfortunately, most of sources in our sample have no well-measured X-ray spectral index, which precludes us to do so. Tentatively in this work, whether or not to include X-ray data in the fit was decided solely based on a visually inspection on the SED of each object. Similar to Wu et al. (2007), the objects were assigned an LBL/IBL/HBL classification according to  $\nu_{\rm peak}$ : for LBLs  $\log \nu_{\rm peak} < 14.5$ , for IBLs  $14.5 < \log \nu_{\rm peak} < 16.5$ , and for HBLs  $\log \nu_{\rm peak} > 16.5$ .

We estimated the jet power using the formula derived from Punsly (2005):

$$Q_{\text{iet}} = 5.7 \times 10^{44} (1 + z)^{1+\alpha} Z^2 F_{151} \text{ erg s}^{-1}$$
 (2)

<sup>1</sup> http://altamira.asu.cas.cz/iblwg/catalog.php

$$Z \approx 3.31 - 3.65 \times [(1+z)^4 - 0.203(1+z)^3 + 0.749(1+z)^2 + 0.444(1+z) + 0.205]^{-0.125}$$
(3)

where  $F_{151}$  is the optically thin flux density from the lobes measured at 151 MHz in units of janskys, and the value of  $\alpha \approx 1$  is a good fiducial value (see Punsly, 2005, for more details). We calculated the  $F_{151}$  through the extrapolation from either the extended radio flux at 5 GHz or 1.4 GHz, or the low frequency total radio flux assuming  $\alpha$ =1.0.

To estimate the black hole mass, we collected the R-band magnitude of host galaxies for 121 sources, of which the core R-band magnitude are also available. Following Wu et al. (2007), we adopted the average redshift of the classification (i.e. LBLs, IBLs and HBLs) for sources without redshift. The sample is listed in Table 1: Col. (1) the source IAU name (J2000), Col. (2) the source alias name, Col. (3) the redshift, Col. (4) the intrinsic synchrotron peak frequency, Col. (5) the observed synchrotron peak luminosity, Col. (6) the jet power, Col. (7) the Doppler factor measured in Wu et al. (2007), Cols. (8) - (9) the R-band magnitude of the host galaxies and the corresponding references, respectively, Col. (10) the black hole mass.

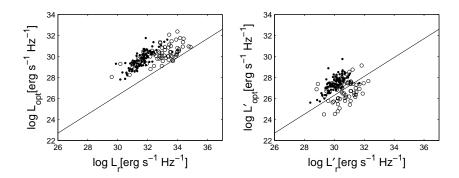
### 3 CORRELATION ANALYSIS

It is commonly accepted that the luminosity of blazars is dominated by the emission from jets, which however is Doppler boosted. Using the Doppler factor derived in Wu et al. (2007) and assuming a same Doppler factor at various wavebands of synchrotron emission as suggested by Capetti et al. (2002) and Trussoni et al. (2003), we calculate the intrinsic luminosity using  $L_{\rm int} = L_{\rm obs}/\delta^{2+\alpha}$  (corresponding to a continuous jet), in which  $L_{\rm int}$  is the intrinsic luminosity,  $L_{\rm obs}$  the observed luminosity. Since the synchrotron peak frequency  $\nu_{\rm peak} \propto B\delta\gamma_{\rm peak}^2$ , the intrinsic one  $\nu_{\rm peak}'$  can be obtained as  $\nu_{\rm peak}/\delta$ , which is only dependent of B and  $\gamma_{\rm peak}$ , where B is the strength of the magnetic field,  $\delta$  the Doppler factor, and  $\gamma_{\rm peak}$  a characteristic electron energy that is determined by a competition between accelerating and cooling processes. With the intrinsic parameters, we explore various correlations below.

## 3.1 Radio and optical core luminosity

We plot the observed radio and optical core luminosity in the left panel of Fig. 1, together with the best fit (the solid line) of the correlation found by Chiaberge et al. (1999) for FR I nuclei. The objects in the present sample span about five orders of magnitude in the radio core luminosity, and the optical core luminosity increase linearly with it over this range. There is a significant correlation between the radio and optical core luminosity at  $\gg 99.99\%$  level, which is confirmed by the partial correlation analysis to exclude the common dependence on the luminosity distance. This correlation suggests that the emission at two wavebands is ascribed to the same non-thermal process, just like in FR I radio galaxies. However, we find that the points are all above the fitted line of Chiaberge et al. (1999), and the optical core luminosity are brighter with about two orders of magnitude than that of FR I nuclei, at the same radio core luminosity.

We correct the beaming effect for the radio and optical core luminosity with the estimated Doppler factor. The intrinsic optical core luminosity are estimated assuming  $\alpha_0 = 1$  (Falomo et al. , 1994), while we use  $\alpha_r = 0.0$  in the radio band. We also recalculated the luminosities in the hypothesis of a bulk Lorentz factor of  $\Gamma = 5$  and a viewing angle of  $60^{\circ}$ , which is the average angle expected for

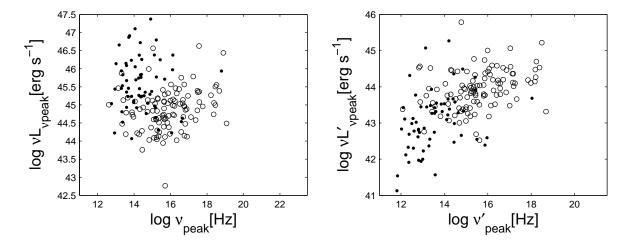


**Fig. 1** The optical vs. radio core luminosity. The solid line is the best fit of the correlation found by Chiaberge et al. (1999) for FR I nuclei. The open circles represent LBLs, while the filled circles are for IBLs and HBLs. In the left panel, the observed luminosities are shown, while the debeamed luminosities with  $\Gamma = 5$  for a mean viewing angle of  $60^{\circ}$  are used in the right panel (see text for details).

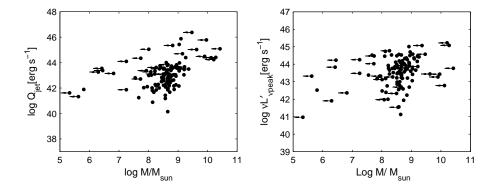
the population of FR I radio galaxies used by Chiaberge et al. (1999) to derive the correlation. The recalculated luminosities are presented in the right panel of Fig. 1. It's apparent that the data points move to the lower radio and optical luminosity due to the debeaming. Although with some scatter, the debeamed radio and optical luminosity follow the correlation of Chiaberge et al. (1999) for FR I nuclei. This supports a common nature of BL Lac objects and FR I radio galaxy nuclei, which is consistent with the results of Giroletti et al. (2006) from the study of 29 nearby BL Lac objects.

# 3.2 Synchrotron peak frequency and luminosity

In the left panel of Fig. 2, we present the relationship between the observed peak luminosity  $\nu L_{\nu peak}$  and synchrotron peak frequency  $\nu_{peak}$ , while the relation of the intrinsic parameters  $\nu L_{\nu_{\rm peak}}$  and  $\nu_{\rm peak}^{'}$  is shown in the right panel. We find a weak negative correlation between  $\nu L_{\nu peak}$  and  $\nu_{peak}$ . In contrast, after correcting the beaming effect both in the peak luminosity and synchrotron peak frequency, we find a significant positive correlation between  $\nu L_{\nu_{\rm peak}}^{'}$  and  $\nu_{\rm peak}^{'}$  with a Spearman correlation coefficient of r = 0.59 at  $\gg 99.99\%$  confidence level. It should be noted that this correlation may be caused by the common dependence of the both parameters on the Doppler factor, as  $L_{\rm int} = L_{\rm obs}/\delta^{2+\alpha}$  ( $\alpha = 1$  in this case) and  $v'_{\text{peak}} = v_{\text{peak}}/\delta$ . We therefore use the partial Spearman correlation method (Macklin 1982) to check this correlation. Still, a significant correlation with a correlation coefficient of 0.33 is present at 99.99% significance level between  $\nu L_{\nu_{\rm peak}}$  and  $\nu_{\rm peak}^{'}$ , independent of the Doppler factor. To further check this correlation, we also perform a statistic analysis on the sources in the restricted Doppler factor range  $\delta$  <3.5. For this subsample of sources, there is no correlation between the Doppler factor and  $v'_{peak}$ , while a significant correlation is still present between  $\nu L_{\nu_{\rm peak}}$  and  $\nu_{\rm peak}$  with a Spearman correlation coefficient of r=0.37, at over 99% significance level (see Fig. 2). Therefore, we conclude that



**Fig. 2** The peak luminosity versus synchrotron peak frequency. The left panel is for observed values, while the intrinsic ones are shown in the right panel. The open circles represent the sources with the Doppler factor  $\delta < 3.5$  (see text for details).



**Fig. 3** Left: the black hole mass versus the jet power. Right: the black hole mass versus the intrinsic luminosity at synchrotron peak frequency. The arrows indicate the upper limit of black hole mass.

the positive correlation between the intrinsic peak luminosity and intrinsic peak frequency might be intrinsic, at least for our present sample.

The intrinsic luminosity at synchrotron peak frequency can be a good indicator of the jet emission due to the fact that the majority of jet synchrotron emission are radiated at the synchrotron peak frequency. Therefore, the significant positive correlation between  $\nu L_{\nu_{\rm peak}}$  and  $\nu_{\rm peak}'$  implies that the jet emission are tightly related with the synchrotron peak frequency, which is only dependent of B and  $\gamma_{\rm peak}$ .

## 3.3 The black hole mass and luminosity

We estimated the black hole mass using the relation between the R-band absolute magnitude of host galaxy  $M_R$  and the black hole mass  $M_{bh}$  proposed by McLure & Dunlop (2002):

$$\log M_{\rm bh} = -0.50 \, M_{\rm R} - 2.96 \tag{4}$$

for 121 sources with available  $M_R$  from literature. There are four sources without measured redshift, which are excluded from the correlation analysis.

Fig. 3 shows the correlation between the black hole mass and jet power (left), and intrinsic peak luminosity (right). The black hole mass of 34 sources are only given as an upper limit due to the faintness and the upper limit of R-band magnitude of host galaxies. We used the astronomy Survival Analysis (ASURV) package (Isobe & Feigelson 1990; Lavalley et al. 1992) when we investigate the correlation involving these upper limits. We find a strong correlation between the intrinsic peak luminosity  $\nu L_{\nu_{\rm peak}}$  and the black hole mass  $M_{\rm bh}$  with a correlation coefficient of r=0.27 at 99.7% confidence level, while there is originally no correlation between the observed peak luminosity and the black hole mass. The partial correlation analysis to exclude the common dependence on redshift still show a significant correlation at 99.9% significance level. Moreover, we find a mild correlation between the jet power and the black hole mass at about 95% confidence level, which is also confirmed by the partial correlation analysis. These results imply that the sources with larger jet power and/or jet emission systematically possess heavier black holes.

## 4 DISCUSSION

It can be seen from Fig. 2 that very high synchrotron peak frequency exist in several objects with  $\log \nu_{\rm peak} > 17$ . Ghisellini et al. (1999) have suggested the existence of ultra-high-energy synchrotron peak BL Lacs (UHBLs), whose synchrotron peak frequencies lies at even higher frequencies than that of conventional HBLs,  $\log \nu_{\rm peak} > 19$ . In our sample, we find eight objects with  $\log \nu_{\rm peak}^{'} > 18$  (see Table 1), which thus can be potential UHBLs. However, as Nieppola et al. (2006) pointed out, using a simple parabolic function in the fitting may produce some error, especially among HBLs, and the peak frequencies of the most extreme objects can be exaggerated and cannot be considered as definite. Although there are no sources with extreme peak frequency  $\log \nu_{\rm peak}^{'} > 19$  in our sample, we conservatively re-examine our results after excluding all eight sources with  $\log \nu_{\rm peak}^{'} > 18$ . We find it does not alter our results.

We note that equation (1) used to estimate the Doppler factor is based on the assumption of a Lorentz factor of five (Giovannini et al., 2001). As Giovannini et al. (2001) claimed, the similar correlation can be obtained for  $\Gamma$  in the range 3-10, so generally equation (1) can be extensively used in this range. However, the simplification of similar Lorentz factor in a sample may produce errors in the estimated Doppler factor. As a matter of fact, the superluminal motion studies (Vermeulen & Cohen , 1994; Jorstad et al. , 2001, 2005) showed that some LBLs may have large apparent velocity, > 10c in some cases, which implies that their Lorentz factor are larger than 10. We collected the available apparent velocity for 27 sources of our sample from literatures, and find that the median value is about 3c, and only three sources with apparent velocity larger than 10c. As BL Lac objects are normally viewed at small viewing angle, the Lorentz factor can be roughly constrained if we assume the viewing angle  $\theta = 1/\Gamma$ . We find that the median value of the constrained Lorentz factor is right in the range 3-10. Moreover, utilizing the variability Doppler factor and apparent velocity,

Lähteenmäki & Valtaoja (1999) and Valtaoja et al. (1999) showed that the averaged lorentz factors of radio selected BL Lac objects are around 5. In addition, Gabuzda et al. (2000) found no evidence that the typical jet Lorentz factor exceed  $\Gamma \sim 5-6$  from the apparent superluminal speed distribution of 16 northern BL Lac objects. We thus conclude that the assumption of a Lorentz factor of 5, or a range of 3-10 is a good approximation although the exceptions do exist. We find that the sources with large apparent velocity are mainly LBLs. If these sources do have Lorentz factor larger than ten, our method may likely underestimate the Doppler factor. Adopting a real Doppler factor may produce a smaller intrinsic synchrotron luminosity than ours, which however may strengthen our results since these sources will be lower down vertically in the LBL region in Fig. 2.

The 'blazar sequence', power -  $v_{peak}$  anti-correlation, was originally presented by Fossati et al. (1998) based on one radio- and one X-ray-selected BL Lacs samples, and one FSRQ sample. These samples had been assembled in an independent and somewhat different way, especially so as regarding the selection band, and indeed none of the individual samples in Fig. 7 of Fossati et al. (1998) shown the claimed anti-correlation between power and  $\nu_{\text{peak}}$ , which was only apparent by combining the three samples (Padovani, 2007). The blazar sequence was recently tested by Nieppola et al. (2008). After eliminating the beaming effect on peak frequencies and luminosities, Nieppola et al. (2008) found that the blazar sequence disappears. Instead, for BL Lac objects separately, they found a positive correlation between the synchrotron peak frequencies and luminosities. Intriguingly, this is consistent with our results. The blazar sequence thus can be an observational phenomenon created by variable Doppler boosting across the synchrotron peak frequency range (Nieppola et al., 2008). By now which really determines the shifting of the synchrotron peak frequency seems still unclear. Wang et al. (2002) claimed that the shape of SEDs is related to the accretion rate in BL Lac objects. They found that HBLs and LBLs follow the different anti-correlation between the peak frequency and accretion rate, which may indicate differences in the physical processes associated with the jets. We re-studied their results by collecting the emission line flux and using the intrinsic peak frequency for our samples. We found that the overall relation between the peak frequency and accretion rate is similar to their results. However, the correlation in HBLs and LBLs separately are not evident. The difference may be partly because we are using the intrinsic peak frequency. On the other hand, the number of BL Lac objects with the detected emission lines is still limited.

For a BL Lac sample, the luminosity at any fixed waveband actually represent different portion of source SED because of the changing  $\nu_{peak}$  in sources. Therefore, the dependence between the source luminosity and synchrotron peak frequency might be necessarily checked at the synchrotron peak frequency, at which the most of jet radiation are emitted. Moreover, it is necessary to debeam the luminosity, which is Doppler boosted due to the relatively small viewing angles of jets. The strong positive correlation between the intrinsic peak luminosity and the intrinsic synchrotron peak frequency implies that the intrinsic source luminosity does depend on the location of the synchrotron peak, although with a large scatter. It seems that the higher peaked BL Lac objects are intrinsically more powerful than the low frequency peaked BL Lac objects in the present sample. Since  $v'_{\text{peak}} \propto \gamma^2_{\text{peak}} B$ , our results show that the power of jet emission increases with the magnetic field and electron Lorentz factor. This however is hard to be explained by the existing models that advocate that the SED of relativistic jets is strongly affected by the (external) radiation field (e.g. Ghisellini et al., 1998). It may be possible that the cooling process is not important in BL Lac ob-

jects. As a matter of fact, the cooling in BL Lac objects is expected to be smaller than that in Flat Spectrum Radio Quasars (FSRQs) which have strong broad emission lines. Alternatively, the efficient injection of the fresh high energy electrons can be maintained for certain long time, although the cooling process can not be ignored. In any cases, the theoretical models should be developed to explain our results. It may also be important if we could measure the magnetic field for individual sources.

We note that the strong-lined counterparts of HBLs, i.e. FSRQs with high  $\nu_{peak}$  have also been discovered in DXRBS by Padovani et al. (2003). Solely with DXRBS, they found no anti-correlation between synchrotron peak frequency and radio, BLR and jet powers, contrary to the predictions of the blazar sequence scenario. Due to the great similarities between FSRQs and BL Lacs, it might be crucial to revisit our results with a large and complete sample including both BL Lacs and FSRQs, which will help us understand the dependence between the source luminosity and synchrotron peak frequency, jet formation and physics.

## **5 CONCLUSIONS**

By eliminating the beaming effect with the estimated Doppler factor from the radio data (Wu et al., 2007), we have calculated the intrinsic luminosity and the intrinsic synchrotron peak frequency for a sample of 170 BL Lac objects. Moreover, we estimated the black hole mass for our sample. The results can be summarized as follows:

- 1. After correcting the beaming effects, the radio and optical core luminosities follow with little scatter the same correlation found for FR I radio galaxies by Chiaberge et al. (1999), in support of the unification of BL Lac objects and FR I galaxies based on orientation.
- 2. We find a significant positive correlation between the intrinsic luminosity at synchrotron peak frequency and the intrinsic synchrotron peak frequency. The result implies that the jet emission is tightly related with the synchrotron peak frequency.
- 3. After correcting the beaming effect, we find a strong correlation between the black hole mass and the peak luminosity, and a mild correlation in the case of jet power. The results indicate that the more powerful jets may have heavier black holes.

Acknowledgements We thank the anonymous referee for insightful comments and constructive suggestions. We thank Xinwu Cao for helpful discussions and suggestions. This work is supported by the NSFC under grants 10373019, 10633010 and 10703009. The National Radio Astronomy Observatory is operated by Associated Universities, Inc., under cooperative agreement with the National Science Foundation. MERLIN is a National Facility operated by the University of Manchester at Jodrell Bank Observatory on behalf of PPARC. This research made use of the NASA/IPAC Extragalactic Database (NED), which is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration. This work also made use of Astrophysical Catalogues Support System (CATS) maintained by the Special Astrophysical Observatory, Russia.

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 Table 1
 The sample

IAU name	Source	Z	Log v' <sub>peak</sub>	$\text{Log } \nu L_{\nu p}$	$\text{Log }Q_{ ext{jet}}$	δ	$m_{ m h}$	Refs.	Log M <sub>bh</sub>
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)
0006-063	NRAO 5	0.347	12.17	45.44	44.79	5.97	•••	•••	
0007 + 472	RX J0007.9+4711	0.280	14.75	44.62	43.37	2.71	>19.01	1	7.93
0035 + 598	1ES 0033+595	0.086	17.01	44.58	41.87	2.33	>17.56	2	7.25
0040 + 408	1ES 0037+405	0.271	18.17	44.99	42.79	1.92			
0050 - 094	PKS 0048-097	0.216	12.89	45.13	44.22	3.29			
0110+418	NPM1G +41.0022	0.096	15.23	44.05	42.70	1.06	15.28	1	8.51
0112 + 227	S2 0109+22	0.473	12.92	46.23	44.17	8.50			
0115 + 253	RXS J0115.7+2519	0.350	13.47	44.80	43.59	2.59	17.91	1	8.76
0123 + 343	1ES 0120+340	0.272	17.99	45.51	42.72	2.93	17.41	3	8.69
0124+093	MS 0122.1+0903	0.339	15.61	44.38	41.43	1.96	18.20	4	8.57
0136+391	B3 0133+388	0.271	16.20	45.75	43.29	2.44	•••		
0141 - 094	PKS 0139-09	0.733	13.19	46.06	45.03	7.43	>18.09	4	9.63
0148 + 140	1ES 0145+138	0.125	15.86	43.94	42.01	0.54	16.67	2	8.12
0153 + 712	8C 0149+710	0.022	14.72	44.03	41.86	1.29	11.68	1	8.65
0201 + 005	MS 0158.5+0019	0.299	17.85	45.27	42.32	2.30	17.84	2	8.59
0208 + 353	MS 0205.7+3509	0.318	15.18	44.73	41.76	2.75			•••
0214 + 517	87GB 02109+5130	0.049	15.51	44.12	42.33	1.50	13.56	3	8.60
0222+430	3C 66A	0.440	14.84	46.25	45.41	3.85			•••
0232+202	1ES 0229+200	0.140	16.48	44.96	43.04	1.93	15.34	2	8.92
0238 + 166	AO 0235+164	0.940	12.68	46.87	45.09	10.78	>17.18	4	10.42
0301 + 346	MS 0257.9+3429	0.245	13.29	44.78	42.34	1.90	17.18	2	8.68
0314+247	RXS J0314.0+2445	0.054	13.36	44.47	41.25	0.97	15.65	1	7.67
0326+024	2E 0323+0214	0.147	15.87	44.72	42.01	2.07	16.59	1	8.36
0416+010	2E 0414+0057	0.287	16.42	45.57	43.98	2.13	16.83	2	9.05
0422 + 198	MS 0419.3+1943	0.512	16.10	45.44	43.01	2.32	18.78	2	8.81
0424+006	PKS 0422+004	0.310	14.88	45.71	44.15	6.24			•••
0505+042	RXS J0505.5+0416	0.027	15.62	42.76	41.89	1.20	17.81	1	5.81
0507+676	1ES 0502+675	0.314	18.04	45.94	42.05	5.62		•••	•••
0508 + 845	S5 0454+84	0.112	12.83	44.07	41.62	10.76	>21.99	2	5.33
0509-040	4U 0506-03	0.304	18.39	45.39	43.42	1.93	17.76	2	8.66
0613+711	MS 0607.9+7108	0.267	14.48	44.84	41.89	3.49	17.00	2	8.87
0625+446	87GB 06216+4441	0.473	13.36	45.69	44.53	5.46			
0650+250	1ES 0647+250	0.203	15.51	45.57	42.71	3.24	>18.61	2	7.73
0654+427	B3 0651+428	0.126	14.98	44.51	42.93	2.38	15.59	1	8.67
0656+426	NPM1G +42.0131	0.059	15.41	44.14	43.73	0.86	13.63	1	8.78
0710+591	EXO 0706.1+5913	0.125	16.85	44.66	42.85	1.87	15.57	1	8.67
0721+713	S5 0716+714	0.300	14.30	45.99	44.35	3.23	>19.55	2	7.74
0738+177	PKS 0735+17	0.424	12.82	46.38	43.28	29.41	>19.75	2	8.08
0744+745	MS 0737.9+7441	0.315	14.52	45.12	42.60	3.05	17.55	2	8.81
0753+538	S4 0749+54	0.200	12.36	44.81	43.26	9.26	>21.44	2 2	6.30
0757+099	PKS 0754+100	0.266	13.00	45.74	43.11	17.72	>18.31		8.21
0806+595	SBS 0802+596	0.300	15.31	45.28	43.04	2.36	16.63	1	9.20
0809+523	1ES 0806+524	0.137	16.14	45.04	42.69	3.32	15.83 >20.47	3	8.65
0818+423 0823+223	OJ 425 4C 22.21	0.530 0.951	12.55	45.53	45.05	6.33	>20.47	2 4	8.01
			12.87	45.88	46.38	2.73		2	9.47
0825+031	PKS 0823+033	0.506	11.84	45.94	43.08	29.42	>19.28	_	8.55
0831+044	PKS 0829+046	0.174	13.07	45.24	43.56	6.21	16.66	2	8.52
0831+087	1H 0827+089 OJ 448	0.941	14.00	45.35	44.95	4.04		2	 9 <i>66</i>
0832+492		0.548	12.38	45.46	43.84	8.39	>19.26		8.66
0854+441 0854+201	US 1889	0.382 0.306	15.04 12.33	45.19 46.09	43.90 43.10	1.79 17.87	 \ 18 08	2	8.50
0854+201 0915+295	OJ 287 B2 0912+29	0.306	12.33	45.32	43.10 44.01	2.96	>18.08		
					43.55		 15.05	 1	 8 08
0916+526	RXS J0916.8+5238	0.190 0.370	15.36 13.33	44.40 45.45		1.85 9.23	15.95	1	8.98
0929+502 0930+498	RXS J0929.2+5013 1ES 0927+500	0.370	17.33	45.45 44.96	43.86 41.85	2.70	 17.36	2	 8.26
0930+498	B2 0927+35	0.188	17.33	44.96 44.52	41.85 44.29	3.68			
0950+550	RGB J0952+656	0.302	12.81 14.96	44.32	43.24	1.99	•••	•••	•••
0954+492	MS 0950.9+4929	0.302	17.25	44.83 45.08	43.24 42.06	2.12	•••	•••	•••
0934+492	S4 0954+65	0.368	17.23	45.08 45.76	43.51	6.79	>18.81	2	8.37
07507055	ひす ひノンサモひン	0.500	13.34	75.70	73.31	0.77	/10.01		0.57

 Table 1
 Continued...

IAU name	Source	z	Log $\nu_{\rm peak}^{'}$	$\text{Log } \nu L_{\nu p}$	Log Q <sub>jet</sub>	δ	$m_{ m h}$	Refs.	Log M <sub>bh</sub>
(1)	(2)	(3)	$\frac{\text{Log } v_{\text{peak}}}{(4)}$	(5)	(6)	(7)	(8)	(9)	$\frac{\text{Log } M_{\text{bh}}}{(10)}$
1012+424	RXS J1012.7+4229	0.364	18.24	45.55	43.69	1.92	17.61	1	8.96
1015+494	GB 1011+496	0.200	16.27	45.14	43.61	2.59	16.15	3	8.94
1031+508	1ES 1028+511	0.360	16.71	45.79	43.10	3.39	18.10	3	8.70
1037+571	RXS J1037.7+5711	0.302	14.40	45.37	43.71	3.70			
1047+546	1ES 1044+549	0.473	12.83	45.46	42.96	2.05	19.31	2	8.44
1053+494	MS 1050.7+4946	0.140	14.99	44.25	42.16	2.49	15.12	1	9.03
1104+382	MRK 421	0.030	16.40	44.67	42.27	1.99	13.22	2	8.23
1109+241	1ES 1106+244	0.460	16.57	45.25	43.78	2.42	18.86	2	8.64
1120+422	EXO 1118.0+4228	0.124	16.33	44.26	42.38	1.76			
1136+701	MRK 180	0.045	16.86	44.15	43.02	0.78	14.37	2	8.10
1136+676 1149+246	RXS J1136.5+6737 EXO 1449.9+2455	0.134 0.402	15.40 17.13	44.50 45.39	41.95 43.60	2.64 2.20	15.76 18.12	3 1	8.66 8.83
1149+240	B2 1147+245	0.402	17.13	43.39	43.15	7.11	>20.40	2	6.82
1150+242	RXS J1151.4+5859	0.302	15.43	45.04	43.79	2.58			
1209+413	B3 1206+416	0.302	13.54	45.22	43.50	6.70			
1215+075	1ES 1212+078	0.136	14.68	44.60	43.13	2.49	15.81	2	8.65
1217+301	B2 1215+30	0.130	14.69	44.59	43.28	3.39	15.78	$\overline{2}$	8.61
1220+345	GB2 1217+348	0.643	13.63	45.84	44.42	5.80		-	
1221+301	PG 1218+304	0.182	17.16	45.36	42.65	2.60	16.86	2	8.47
1221+282	ON 231	0.102	14.45	44.65	42.99	5.32	16.43	2	8.01
1223+806	S5 1221+80	0.473	14.23	45.29	44.77	4.32		•••	
1224+246	MS 1221.8+2452	0.218	13.43	44.73	42.11	2.96	18.32	2	7.96
1230+253	RXS J1230.2+2517	0.135	14.11	44.76	43.08	3.60			
1231+642	MS 1229.2+6430	0.163	15.19	44.60	42.18	2.96	16.15	2	8.70
1237+629	MS 1235.4+6315 1ES 1239+069	0.297 0.150	15.07 18.68	44.70	42.24 41.33	3.02	>22.08	2	 5.62
1241+066 1248+583	PG 1246+586	0.130	18.08	44.49 46.80	41.33 44.70	2.45 11.10	>22.08	$\overset{2}{2}$	5.63 9.15
1253+530	S4 1250+53	0.302	14.17	44.92	44.26	3.52			
1257+242	1ES 1255+244	0.302	16.70	44.86	42.29	1.30	 16.53	2	8.33
1310+325	AUCVn	0.996	12.61	47.10	44.38	27.73	>17.93	4	10.12
1322+081	1ES 1320+084N	1.500	13.14	46.91	44.41	4.09	>18.71	2	10.28
1341+399	RXS J1341.0+3959	0.163	18.21	44.74	43.05	1.44	15.82	1	8.86
1402+159	MC 1400+162	0.244	16.24	45.03	44.40	1.90			
1404+040	MS 1402.3+0416	0.344	15.36	45.42	43.62	1.64	>18.87	2	8.26
1409+596	MS 1407.9+5954	0.496	15.81	44.97	43.35	2.55	18.25	2	9.04
1415+485	RGB J1415+485	0.496	14.16	45.05	43.50	4.06	19.30	1	8.51
1415+133	PKS 1413+135	0.247	12.08	45.00	44.52	3.47		•••	
1417+257	2E 1415+2557	0.237	17.17	45.35	43.57	2.13	16.28	3	9.09
1419+543	OQ 530	0.152	12.93	45.11	42.53	11.38	15.90	2	8.74
1427+238 1427+541	PKS 1424+240 RGB J1427+541	0.160 0.106	15.22 14.13	45.10 44.18	43.55 42.17	2.66 1.38	>20.67 14.81	2 1	6.42 8.86
1427+341	H 1426+428	0.106	18.19	44.18 45.06	42.17	1.56	15.97	2	8.51
1439+395	PG 1437+398	0.125	16.17	45.32	43.83	1.87	17.48	$\frac{2}{2}$	8.95
1442+120	1ES 1440+122	0.163	15.81	44.74	43.10	2.06	16.45	$\frac{2}{2}$	8.55
1444+636	MS 1443.5+6349	0.298	16.58	44.48	42.17	1.69			
1448+361	RXS J1448.0+3608	0.271	14.90	45.16	43.31	2.51			•••
1458+373	B3 1456+375	0.333	12.85	44.93	43.76	5.39		•••	
1501+226	MS 1458.8+2249	0.235	14.46	44.97	42.59	4.59	17.40	2	8.52
1509+559	SBS 1508+561	2.025	14.21	47.38	44.61	5.03			
1516+293	RXS J1516.7+2918	0.130	14.82	44.37	42.87	1.29	15.52	1	8.75
1517+654	1H 1515+660	0.702	17.04	46.63	43.50	3.32	>18.51	2	9.36
1532+302	RXS J1532.0+3016	0.064	15.01	44.10	42.17	1.66	14.86	1	8.26
1533+342	RXS J1533.4+3416	0.810	15.42	45.97	43.55	4.86	 16 07		
1534+372	RGB J1534+372	0.143	13.82	44.22 46.44	41.81	2.21	16.97	1	8.13
1535+533 1536+016	1ES 1533+535 MS 1534.2+0148	0.890 0.312	18.49 17.19	46.44 44.85	44.24 43.39	2.55 1.89	>17.44 17.62	4 2	10.21 8.76
1540+819	1ES 1544+820	0.312	15.63	45.12	43.39	2.30	>17.02	$\frac{2}{2}$	7.82
1540+147	4C 14.6	0.605	13.54	45.83	45.43	6.34	18.76	1	9.04
1542+614	RXS J1542.9+6129	0.302	13.90	45.25	43.06	4.42			

Table 1 Continued...

IAU name	Source	z	Log $\nu_{\rm peak}^{'}$	$\text{Log } \nu L_{\nu p}$	$\operatorname{Log} Q_{\operatorname{jet}}$	δ	$m_{ m h}$	Refs.	Log M <sub>bh</sub>
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)
1554+201	MS 1552.1+2020	0.222	15.81	44.99	42.94	2.06	16.44	1	8.93
1555 + 111	PG 1553+11	0.360	15.42	46.38	44.10	5.07	>20.99	2	7.25
1602+308	RXS J1602.2+3050	0.302	15.34	44.64	43.64	1.85			
1626+352	RXS J1626.4+3513	0.497	15.24	45.08	43.67	2.53	>17.89	1	9.22
1644+457	RXS J1644.2+4546	0.225	15.44	44.63	43.52	1.95	16.74	1	8.79
1652+403	RGB J1652+403	0.240	14.05	44.88	42.89	1.85			
1653+397	MRK 501	0.034	15.96	44.63	41.68	4.79	12.51	1	8.72
1704+716	RXS J1704.8+7138	0.350	14.76	45.04	42.84	2.45	18.52	1	8.45
1719+177	PKS 1717+177	0.137	12.22	44.22	43.05	5.03			
1724 + 400	B2 1722+40	1.049	14.44	46.38	45.66	4.72	•••		
1725+118	H 1722+119	0.170	15.86	45.46	43.38	2.57	>20.75	2	6.45
1728 + 502	IZw187	0.055	16.05	44.02	42.03	1.91	15.36	2	7.84
1739+476	OT 465	0.473	12.02	45.03	44.93	5.42	>19.75	2	8.22
1742+597	RGBJ 1742+597	0.400	13.85	45.24	43.51	3.73	18.67	1	8.55
1743+195	NPM1G +19.0510	0.084	15.46	44.20	42.26	3.19	14.14	1	8.93
1745+398	B3 1743+398B	0.267	16.01	45.11	44.40	1.69	16.22	1	9.26
1747+469	B3 1746+470	1.484	12.50	46.15	45.10	11.33			
1748+700	S4 1749+70	0.770	13.42	46.26	44.49	9.94	>17.68	2	9.90
1749+433	B3 1747+433	0.473	13.61	45.26	44.48	4.46	•••	•••	•••
1750+470	RXS J1750.0+4700	0.160	17.05	44.12	43.14	0.72	16.33	1	8.58
1751+096	PKS 1749+096	0.322	11.81	45.83	42.78	37.09	17.93	2	8.64
1756+553	RXS J1756.2+5522	0.271	16.47	44.97	42.50	1.84		•••	
1757+705	MS 1757.7+7034	0.407	13.17	45.20	42.51	3.02	18.93	2	8.44
1800+784	S5 1803+784	0.680	13.41	46.72	45.35	8.51	>19.51	2	8.82
1806+698	3C 371	0.051	14.18	44.28	43.75	1.79	14.07	1	8.39
1808+468	RGB J1808+468	0.450	14.29	45.14	43.63	2.80	19.22	1	8.42
1811+442	RGB J1811+442	0.350	15.99	44.70	43.70	0.79	17.58	1	8.92
1813+317	B2 1811+31	0.117	14.85	44.09	42.81	1.73	17.66	1	7.55
1824+568	4C 56.27	0.664	12.40	46.14	45.88	4.07	18.88	2	9.10
1829+540	RXS J1829.4+5402	0.302	14.61	44.81	43.51	1.34			
1838+480	RXS J1838.7+4802	0.300	13.65	45.09	43.33	2.45	18.61	1	8.22
1841+591	RGB J1841+591	0.530	14.34	44.93	43.46	1.44	18.08	1	9.21
1853+672	1ES 1853+671	0.212	13.07	44.19	41.67	3.00	17.83	2	8.18
1927+612	S4 1926+61	0.473	12.70	45.67	44.43	7.66			
1959+651	1ES 1959+650	0.047	15.87	44.54	40.89	5.20	14.40	2	8.14
2005+778	S5 2007+77	0.342	12.80	45.46	44.00	7.70	18.16	2	8.60
2009+724 2022+761	S5 2010+72	0.473	12.79 14.49	45.23	44.80 44.43	7.45 5.39	•••	•••	•••
2022+761 2039+523	S5 2023+76 1ES 2037+521	0.473 0.053	14.49 14.78	45.77 44.31	44.43	3.55	 13.65	2	 8 65
2039+323	PKS 2131-021	1.285	14.78	44.31 46.66	40.15 45.78	3.33 11.93	>18.94	$\overset{2}{2}$	8.65 9.96
2134-018 2145+073	MS 2143.4+0704	0.237	13.56	46.66 44.67	43.78 42.99	2.53	>18.94 17.42	$\overset{2}{2}$	9.96 8.52
2143+073 2152+175	PKS 2149+17	0.237	12.57	45.35	42.99	2.33 10.39	>20.63	$\overset{2}{2}$	8.32 7.78
2202+422	BL LAC	0.473	13.58	45.05	42.08	10.39	14.42	2	8.58
2250+384	B3 2247+381	0.070	15.36	43.03 44.64	42.08	2.03	15.51	1	8.64
2257+077	PKS 2254+074	0.119	13.03	45.07	42.51	8.85	16.22	2	8.85
2319+161	O J2319+161	0.190	15.34	44.90	43.41	2.00			
2322+346	TEX 2320+343	0.302	15.18	44.30	42.41	1.23	14.88	 1	8.73
2322+340	1ES 2321+419	0.059	14.21	43.76	41.39	1.70			0.73
2323+421	1ES 2321+419 1ES 2326+174	0.039	16.57	44.89	42.92	2.04	17.17	2	8.51
2339+055	MS 2336.5+0517	0.740	14.78	46.57	44.01	1.82			
2347+517	1ES 2344+514	0.740	15.50	44.23	41.18	3.63	13.39	2	8.57
2350+196	MS 2347.4+1924	0.515	15.88	45.15	43.07	1.81	13.37		0.57
4330T170	1410 2047.471924	0.313	15.00	₹3.13	45.07	1.01	•••	•••	

Notes: Col. 1: the source IAU name (J2000). Col. 2: the source alias name. Col. 3: the redshift. Col. 4: the intrinsic synchrotron peak frequency. Col. 5 the observed peak luminosity in units of erg s<sup>-1</sup>. Col. 6 the jet power in units of erg s<sup>-1</sup>. Col. 7 the Doppler factor. Col. 8 the R magitude of host galaxies. '>' indicates the lower limit. Col. 9 the references of  $m_h$ : 1, Nilsson et al. (2003). 2, Urry et al. (2000). 3, Nilsson et al. (2007). 4, O'Dowd & Urry (2005). Col. 10 the black hole mass in units of solar mass.